

# Energy Component Analysis for Electronically Excited States of Molecules: Why the Lowest Excited State Is Not Always the HOMO/ LUMO Transition

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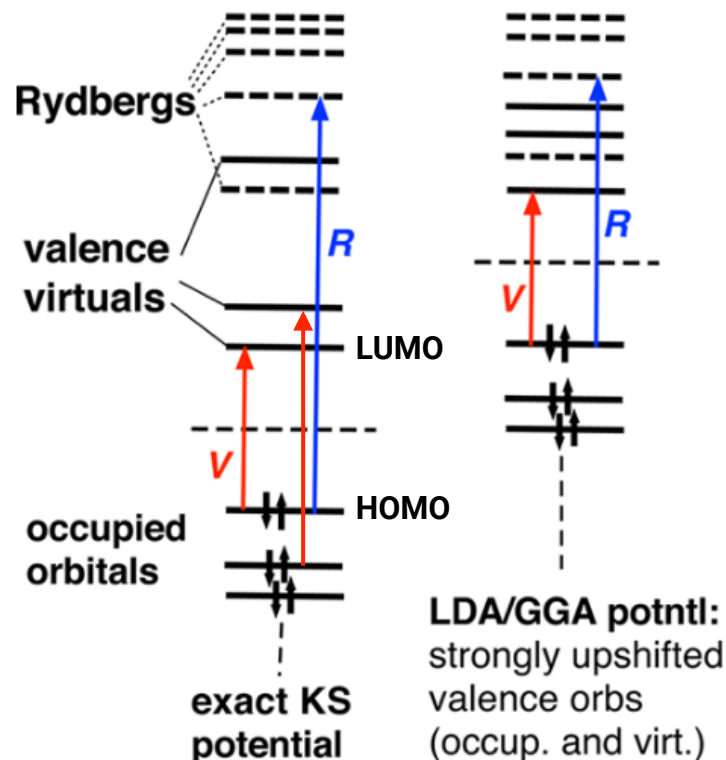
## Theory & Methods

- Practical interpretation of  $K_2$  and  $J_2$
  - Excited-state energy decomposition
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## Results & Discussion

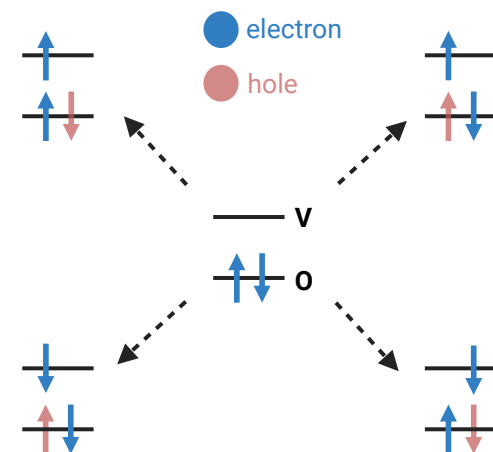
- ACRFLCN: effect of CT and LE states
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## Conclusion



## Two-orbital two-electron model (TOTEM)

Spin-flip excitations      Spin-preserving excitations



- The HOMO and LUMO energies can be strongly method dependent. → **NOT** good for analysis!
- Need to solve many-electron wavefunction. → **Too complicated!**



- Uses transition density matrix (TDM) → **focuses on electron excitation between two orbitals.**
- Reduce to be a two-body wavefunction (electron and hole quasiparticles) → can **captures correlation effects.**



- Individual orbital do not have physical meaning.
- Each excited state is NOT defined by just one pair of orbital.



Optimal representation by natural transition orbitals (NTOs)

The compact orbital obtained from TDM.

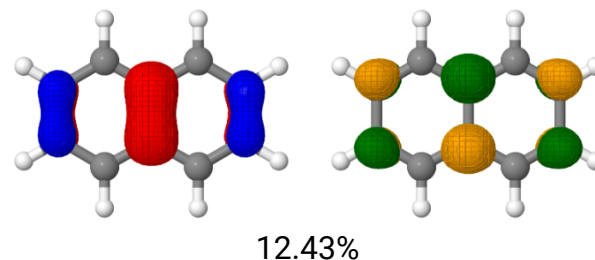
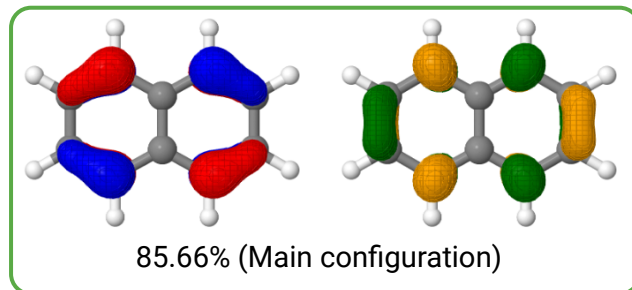
- Compact (few orbital pairs).
- Clear hole–electron picture.

Hole NTO

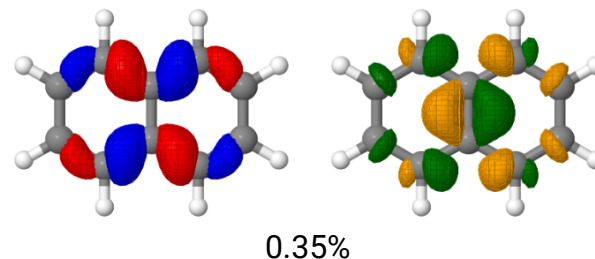
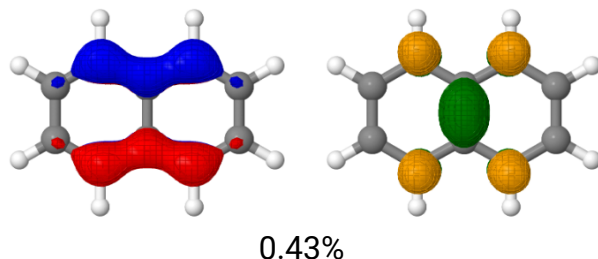
Electron NTO

Hole NTO

Electron NTO

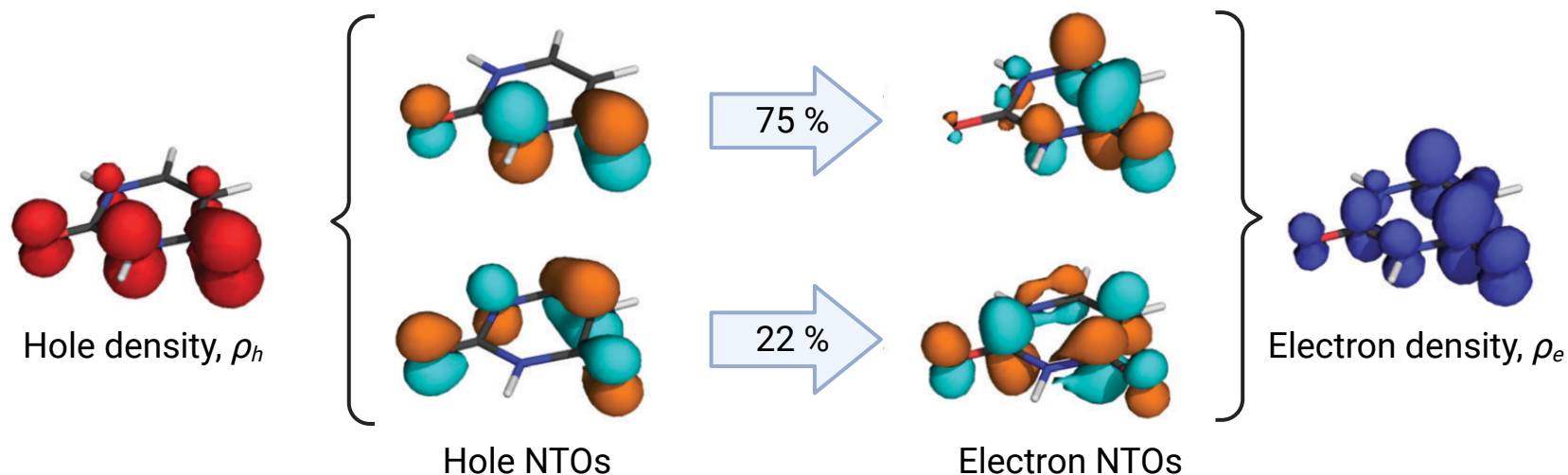


**S<sub>2</sub> state**



Multiple NTO pairs may be needed to fully describe an excited state.

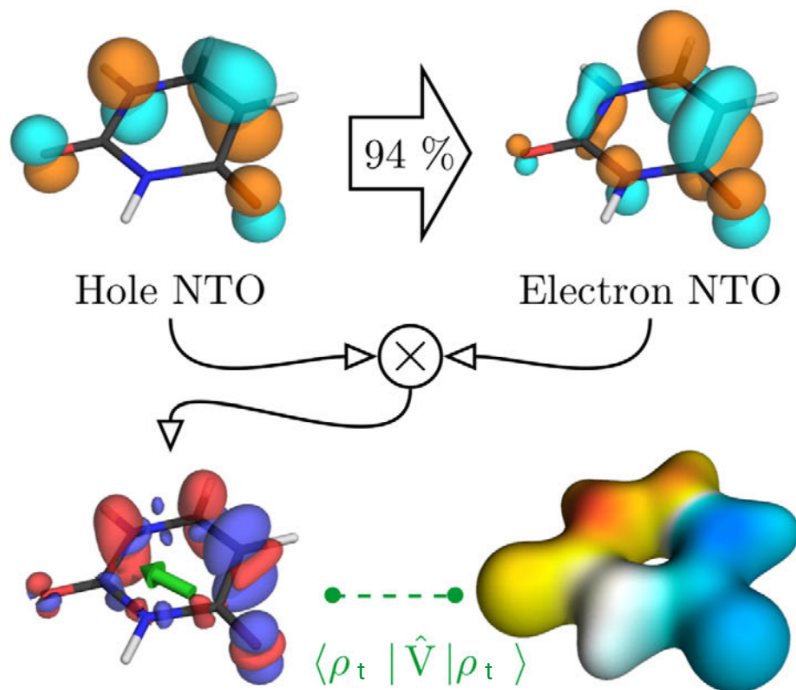
Hole and electron densities represent the overall distribution of hole and electron.



- The NTO pairs can be combined (weighted and summed) to give overall hole and electron densities.
- Represents the area where excited electron come from ( $\rho_h$ ) and where its goes to ( $\rho_e$ ).
- Can be used to represents the **screened coulomb interaction  $K_2$** .

Alternative option: transition density represents the overlap between hole and electron NTOs.

☹️ Shape of the transition density is more difficult to interpret.



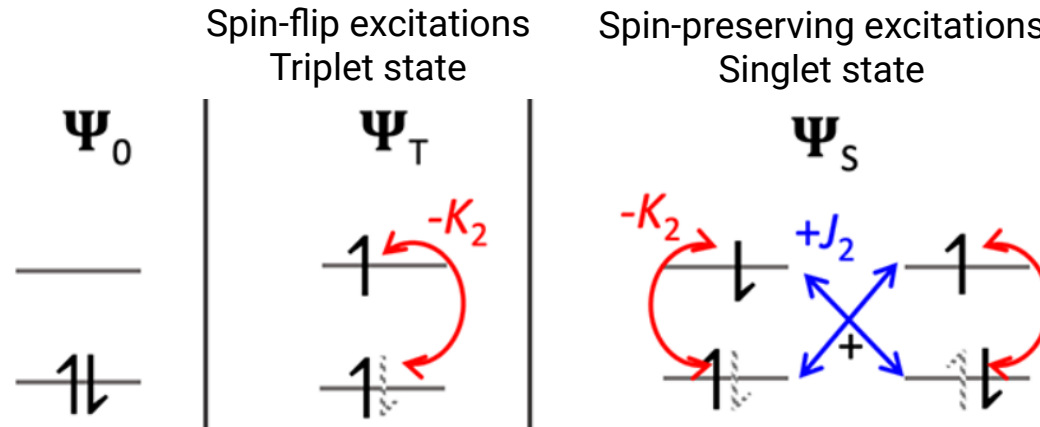
**Transition dipole moment**

$$\vec{\mu}_t = \int \rho_t(r) \vec{x} dr$$

**Oscillator strength**

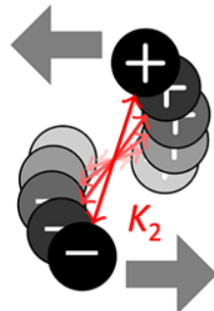
$$f = \frac{2\Delta E}{3} |\vec{\mu}_t|^2$$

- Indicates how strongly an electronic transition interacts with light, measured by its transition dipole moment  $\vec{\mu}_t$  and oscillator strength  $f$ .
- Represents the **transition density repulsion**,  $J_2$ , which is overlap between  $\rho_t$  and their ESPs.



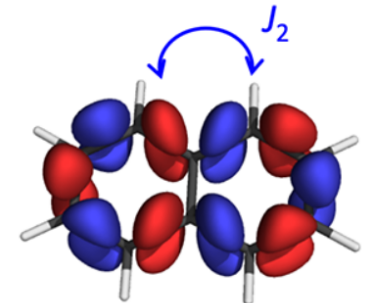
## $K_2$

- Excited electron is attracted by the hole it left behind.
- Applies to both triplet and single states.
- Negative value (stabilize)
- Screened Coulomb interaction between electron and hole.



## $J_2$

- Exchange interaction between excited electron and hole.
- Only applies to singlet state.
- Positive value (destabilize)
- Responsible for the  $S - T$  energy splitting.
- Transition density repulsion.



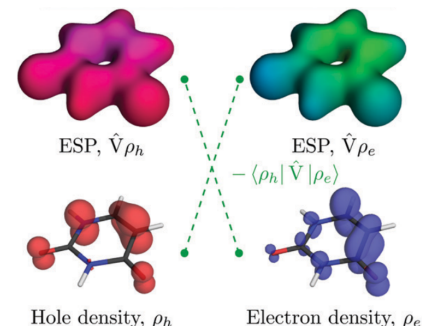
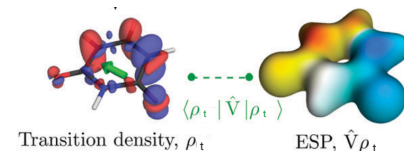
Excitation energies can be calculated as:

$$\Delta E = \underbrace{\sum_{ia} |C_{ia}|^2 (\epsilon_a - \epsilon_i)}_{h'} + J_2 + K_2 + XC_2$$

where,

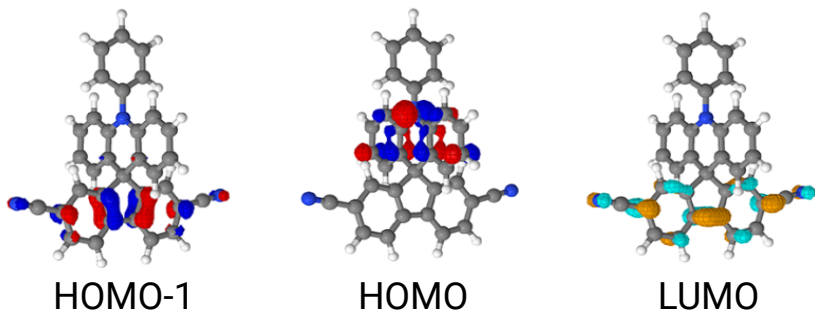
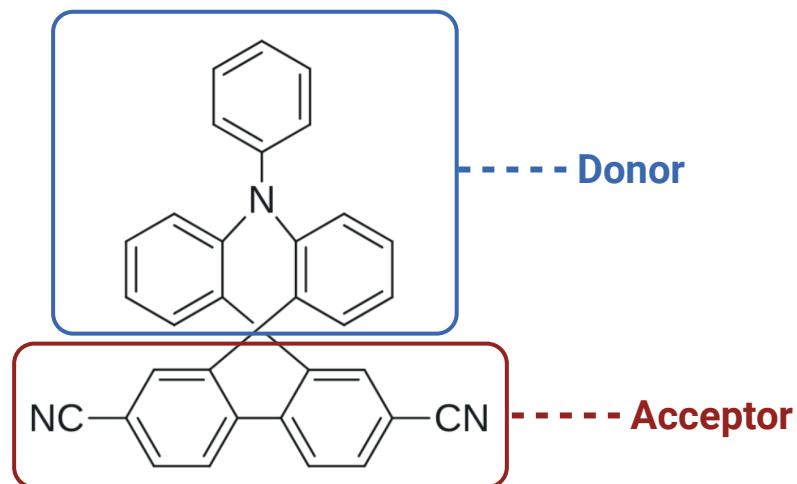
$$J_2 = \iint \frac{\rho_t(r_1)\rho_t(r_2)}{r_{12}} dr_1 dr_2 = \langle \rho_t | \hat{V} | \rho_t \rangle$$

$$K_2 \approx - \iint \rho_h(r_h)\rho_e(r_e) \frac{s(r_{he})}{r_{he}} dr_h dr_e = \langle \rho_h | \hat{V} | \rho_e \rangle$$



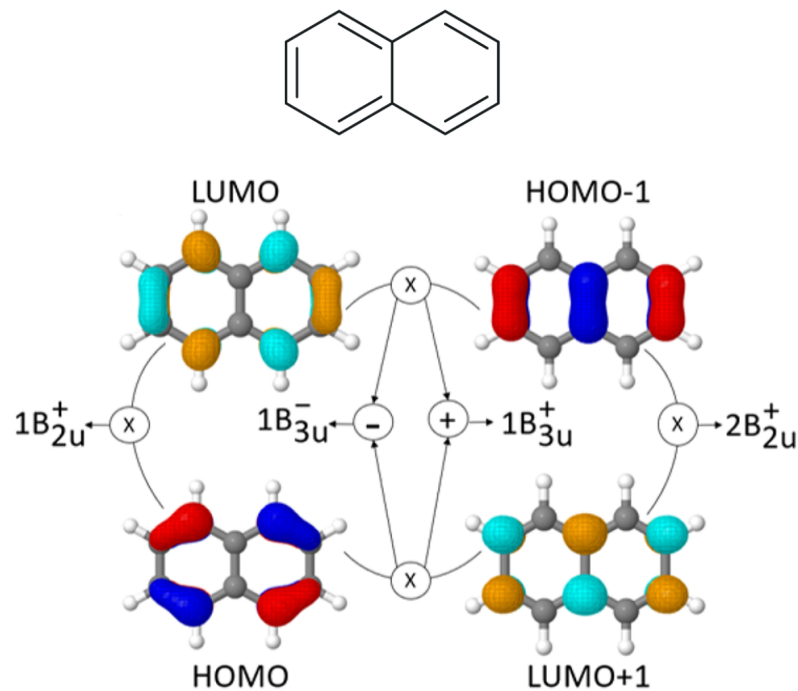
- $h'$  is weighted sum of orbital energy differences.
- $J_2$  is the overlap of the transition density with its own electrostatic potential (ESP).
- $K_2$  is the overlap of the hole density with the ESP from electron density, or *vice versa*.
- $XC_2$  is the remaining part that does not come from the orbital gap, Coulomb attraction, or exchange repulsion (usually small).

**ACRFLCN:** locally excited (LE) VS charge transfer (CT)



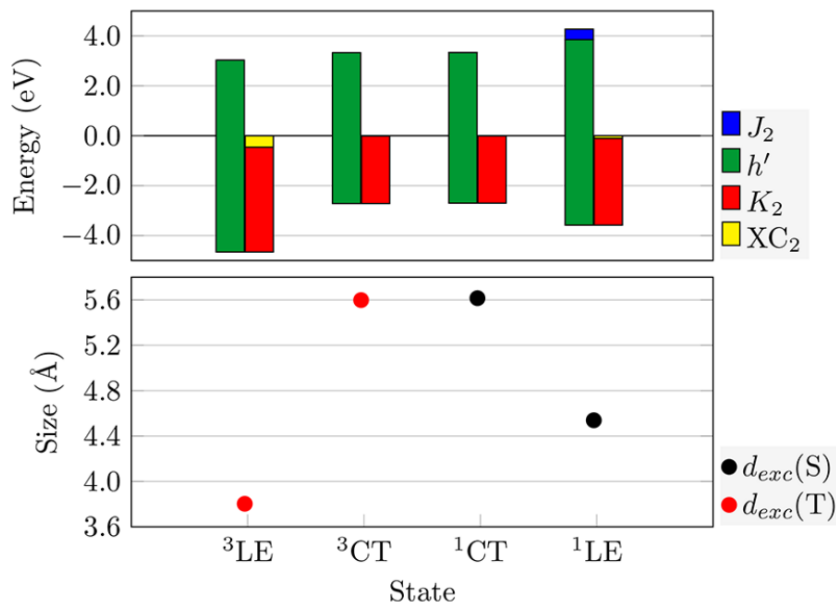
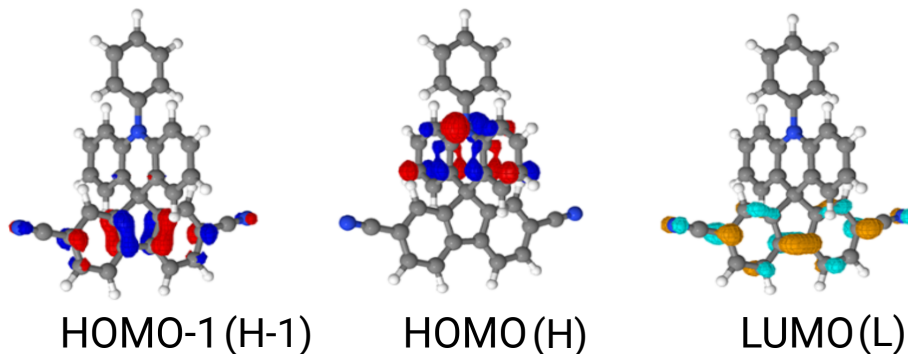
- Contains a separate donor-acceptor system.
- Charge transfer (HOMO → LUMO) VS locally excited (HOMO-1 → LUMO).

**Naphthalene:** ionic "+" VS covalent "-" states



- All low-lying excited state are constructed from the same set of MOs → need additional perspectives to characterize.
- Interactions between configurations affect the oscillator strengths.
- "+" → ionic → bright state → large  $f$ .
- "-" → covalent → dark state → small or zero  $f$ .

## Frontier orbitals



## Excited energy components

$$T_1 (^3LE) = 3.03 \text{ eV} = 7.69 + (-4.20) + (-0.46)$$

$$T_2 (^3CT) = 3.32 \text{ eV} = 6.04 + (-2.71) + (-0.01)$$

$$S_1 (^1CT) = 3.34 \text{ eV} = 6.03 + 0.01 + (-2.69) + (-0.01)$$

$$S_2 (^1LE) = 4.26 \text{ eV} = 7.43 + 0.43 + (-3.48) + (-0.10)$$

## Main Configuration

H-1  $\rightarrow$  L

H  $\rightarrow$  L

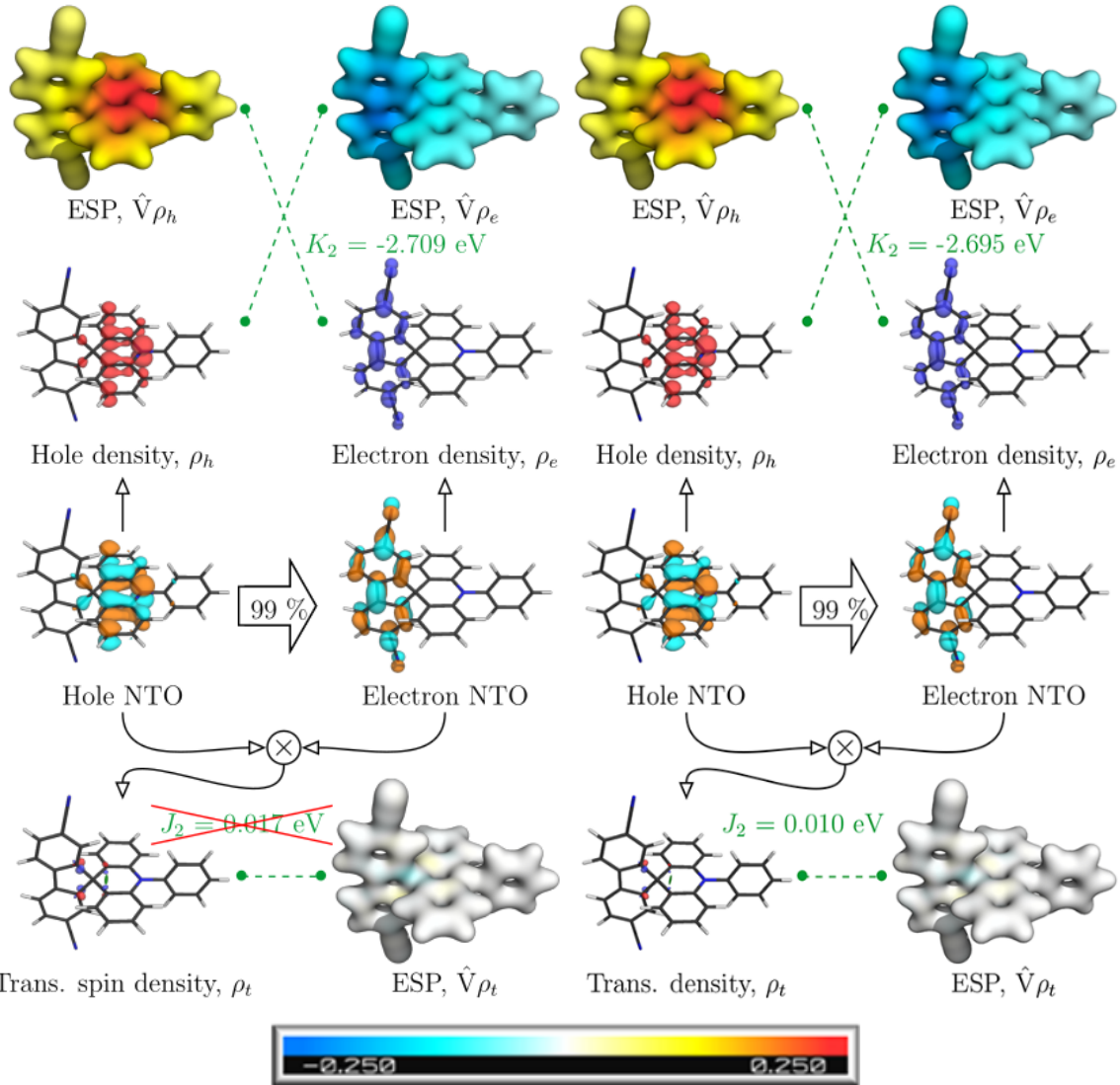
H  $\rightarrow$  L

H-1  $\rightarrow$  L

- $h'$  directly depends on the MO energies.
- $K_2$  term in LE is significantly larger than CT state (more overlap).
- $J_2$  almost vanishes for a state that is completely CT.
- LE vs CT: Exciton size ( $d_{exc}$ ) reflects their difference.

$T_2$  ( $^3CT$ )

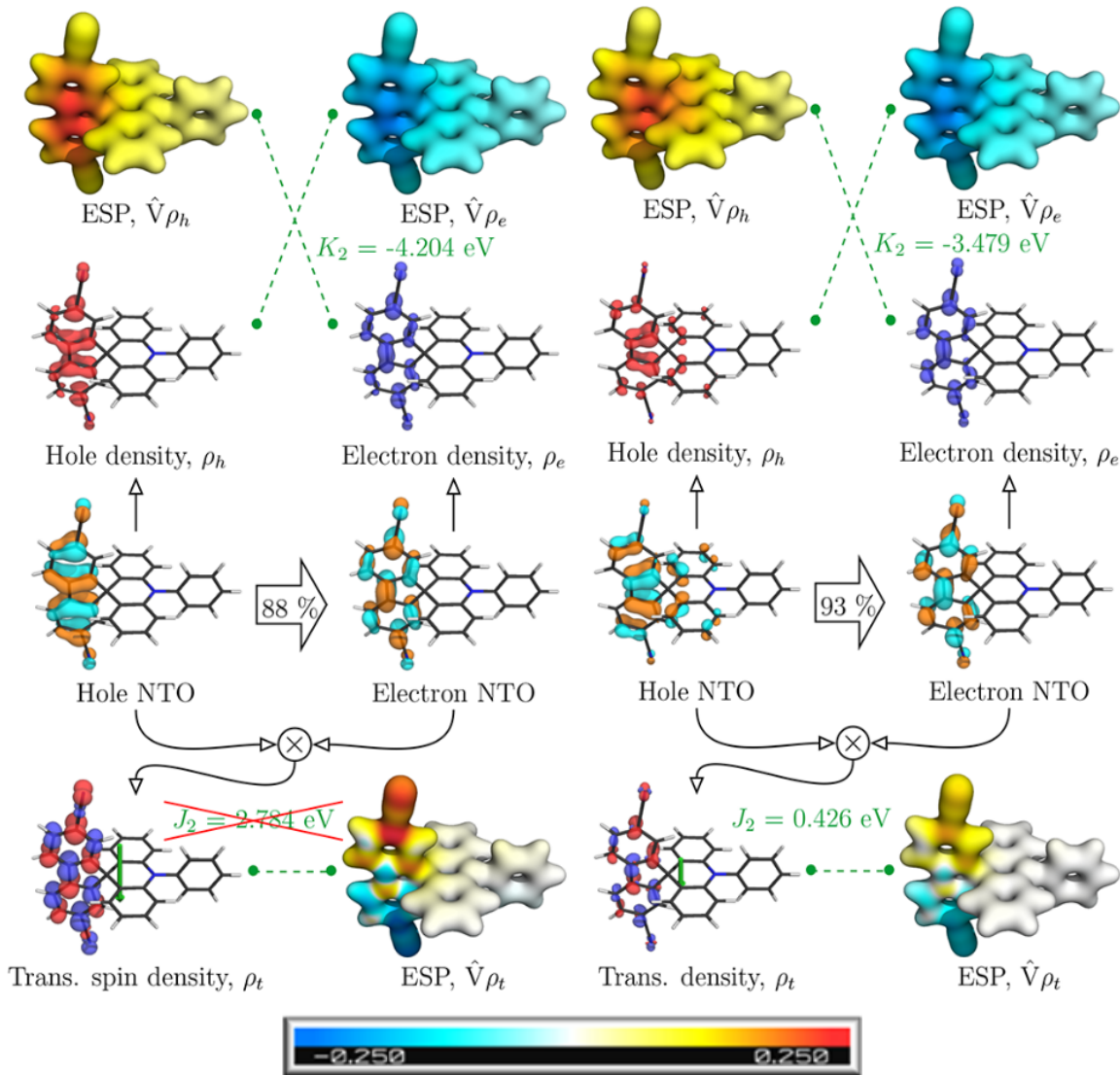
$S_1$  ( $^1CT$ )



- Similar energy contributions and exciton size.
- Only one NTO pair dominates each state → the densities closely match the NTOs.
- HOMO → LUMO (CT)
- $\rho_e$  locates in a region of weaker hole ESP → small  $K_2$  term
- The product between hole and electron NTOs → transition density.
- Transition densities and their ESP almost vanish → very small  $J_2$

$T_1$  ( $^3LE$ )

$S_2$  ( $^1LE$ )



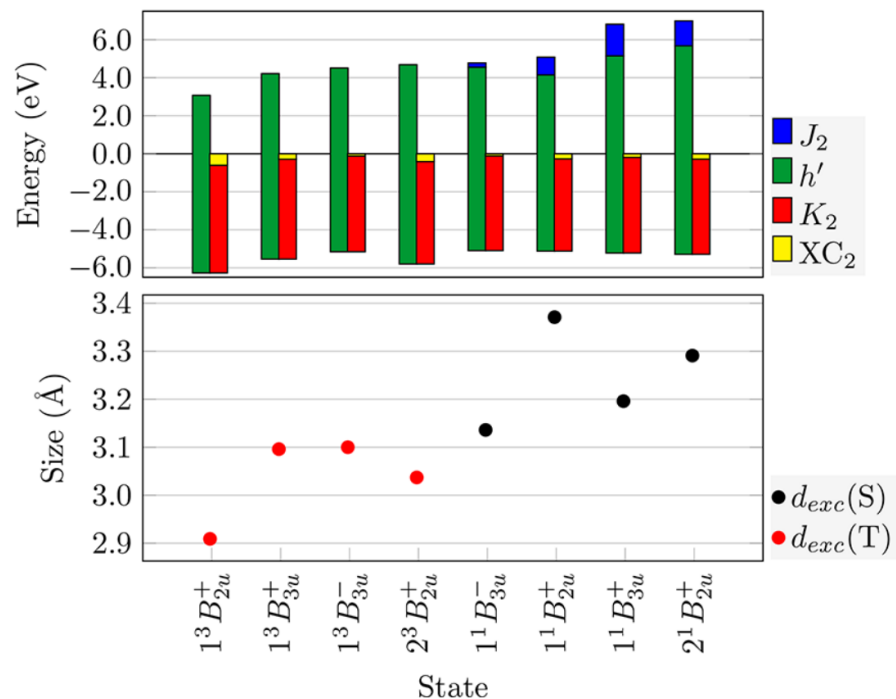
"formal"  $J_2$  for the singlet using the triplet wave functions = 2.78 eV, whereas actual  $J_2$  is only 0.43 eV. → they have different wave functions!



$S_2$  wave functions attempt to lower their energy by minimizing the  $J_2$  → interacts more strongly with other electronic configurations in donor part.



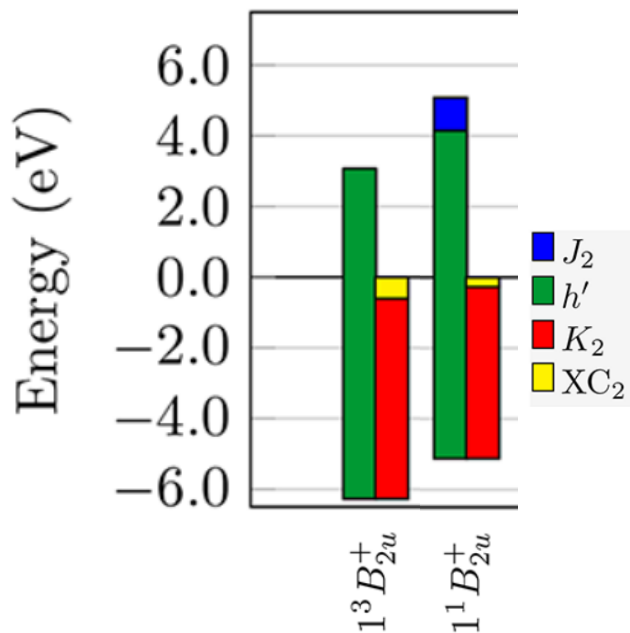
- $d_{exc}$  increase (3.8 → 4.54 Å)
- $K_2$  decrease (-4.20 → -3.48 eV)



State	Main Configuration
$T_1 (1^3B_{2u}^+)$	H → L
$T_2 (1^3B_{3u}^+)$	H-1 → L  +  H → L+1
$T_3 (1^3B_{3u}^-)$	H → L+1  -  H-1 → L
$T_4 (2^3B_{2u}^+)$	H-1 → L+1
$S_1 (1^1B_{3u}^-)$	H → L+1  -  H-1 → L
$S_2 (1^1B_{2u}^+)$	H → L
$S_3 (1^1B_{3u}^+)$	H-1 → L  +  H → L+1
$S_4 (2^1B_{2u}^+)$	H-1 → L+1

- One NTO pair dominates in  $T_1$ ,  $T_4$ ,  $S_2$ , and  $S_4$  states.
- Strong mixing of H → L+1 and H-1 → L transitions forms two states ( $B_{3u}^+$  and  $B_{3u}^-$ ), differing only the sign of their linear combination.
- All excited states are LE ( $d_{exc} < 3.5 \text{ \AA}$ ) → similar  $K_2$ .
- The H-1 → L+1 dominated state is the highest state in both triplet and singlet states ( $T_4$ , and  $S_4$ ).
- $S_1$  is almost identical to  $T_3$ , differing mainly by a presence of  $J_2$  term (0.234 eV).

## $T_1 (1^3B_{2u}^+)$ VS $S_2 (1^1B_{2u}^+)$



Same as ACRFLCN, the "formal"  $J_2$  for H  $\rightarrow$  L singlet state using  $T_1$  wave function is 4.20 eV, while the actual  $J_2$  is 0.94 eV.



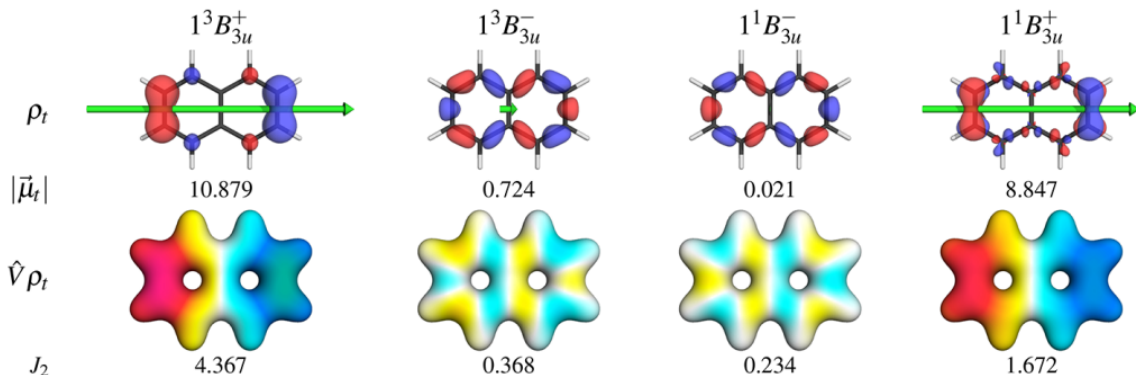
$S_2$  attempt to lower its energy by minimizing the  $J_2$  at the cost of decrease  $K_2$  and  $XC_2$ .

### Excited energy components

$$T_1 (1^3B_{2u}^+) = 3.07 \text{ eV} = 9.35 + (-5.67) + (-0.61)$$

$$S_2 (1^1B_{2u}^+) = 5.09 \text{ eV} = 9.27 + 0.94 + (-4.85) + (-0.27)$$

$B_{3u}$  states (mainly mixtures of H-1  $\rightarrow$  L and H  $\rightarrow$  L+1)



The transition densities of all “+” states are concentrated on the atoms, while in “-” states they are located on the bonds.

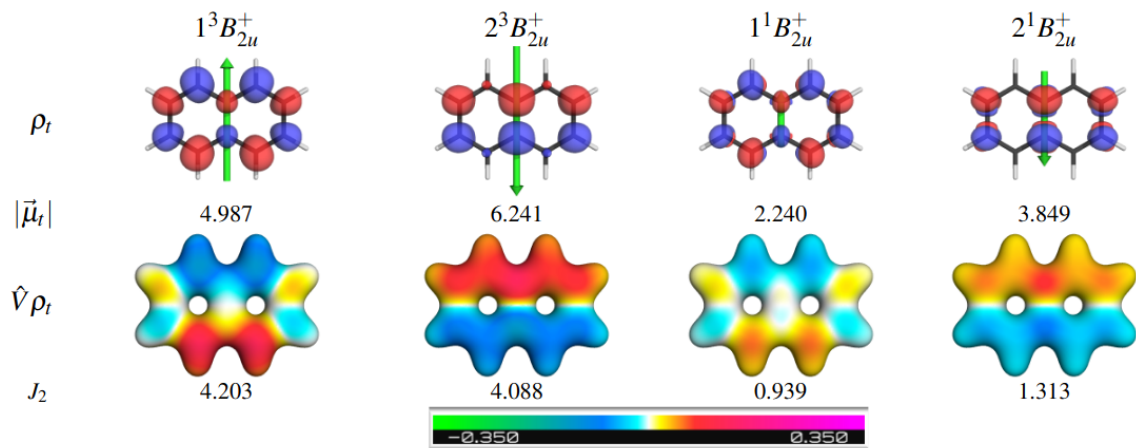


The “+” states show significantly larger transition moments and ESPs than the “-” states.



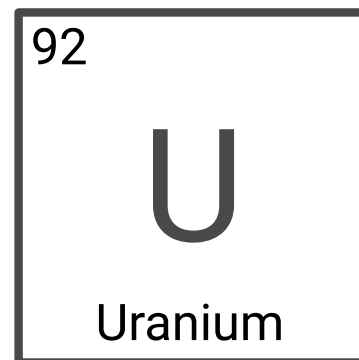
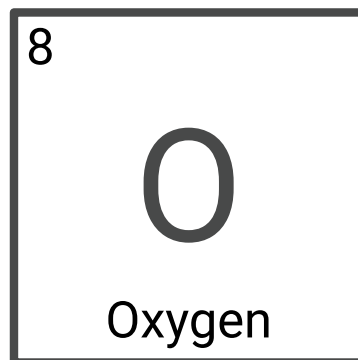
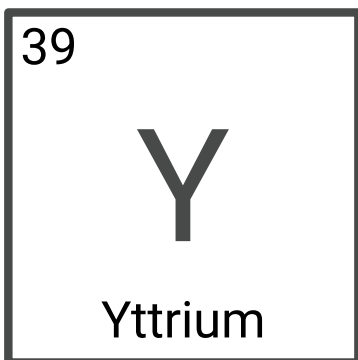
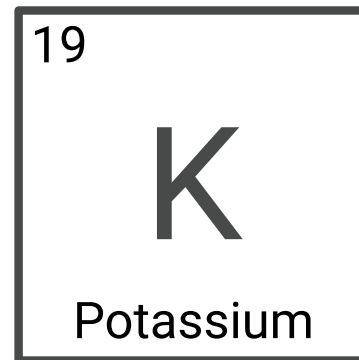
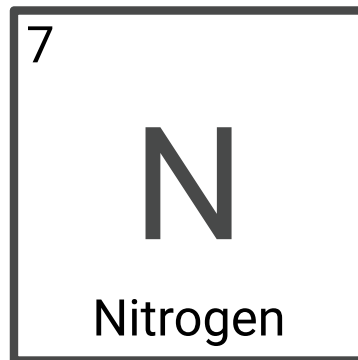
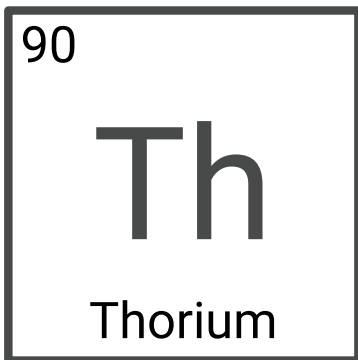
The positive correlation between  $J_2$  and transition moments  $\rightarrow$  brighter singlet states are shifted to higher energy.

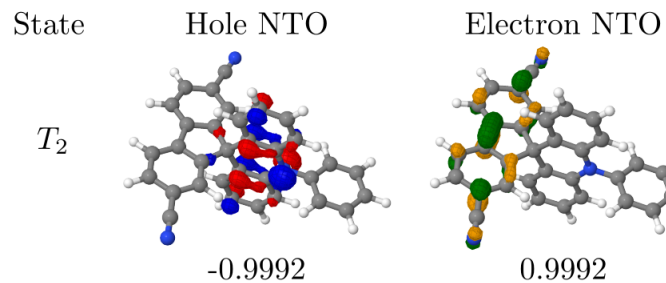
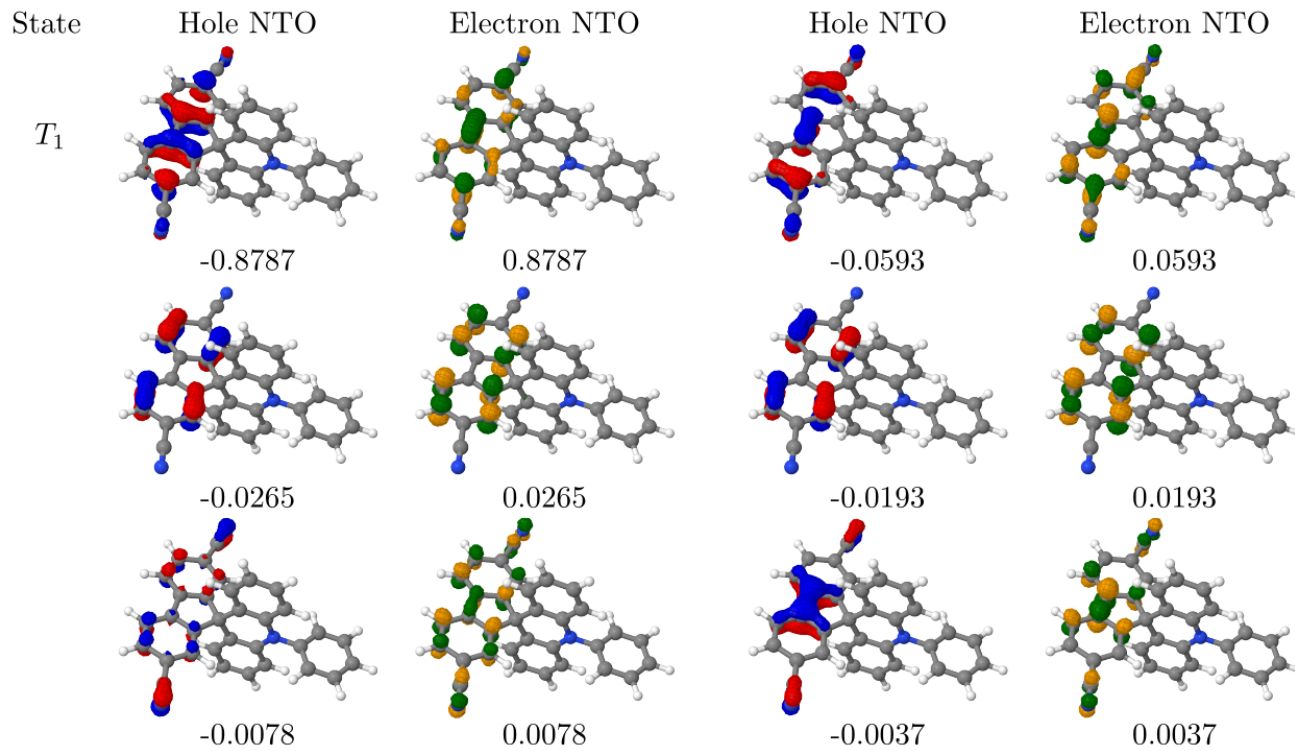
$B_{2u}$  states (mainly H  $\rightarrow$  L or H-1  $\rightarrow$  L+1)

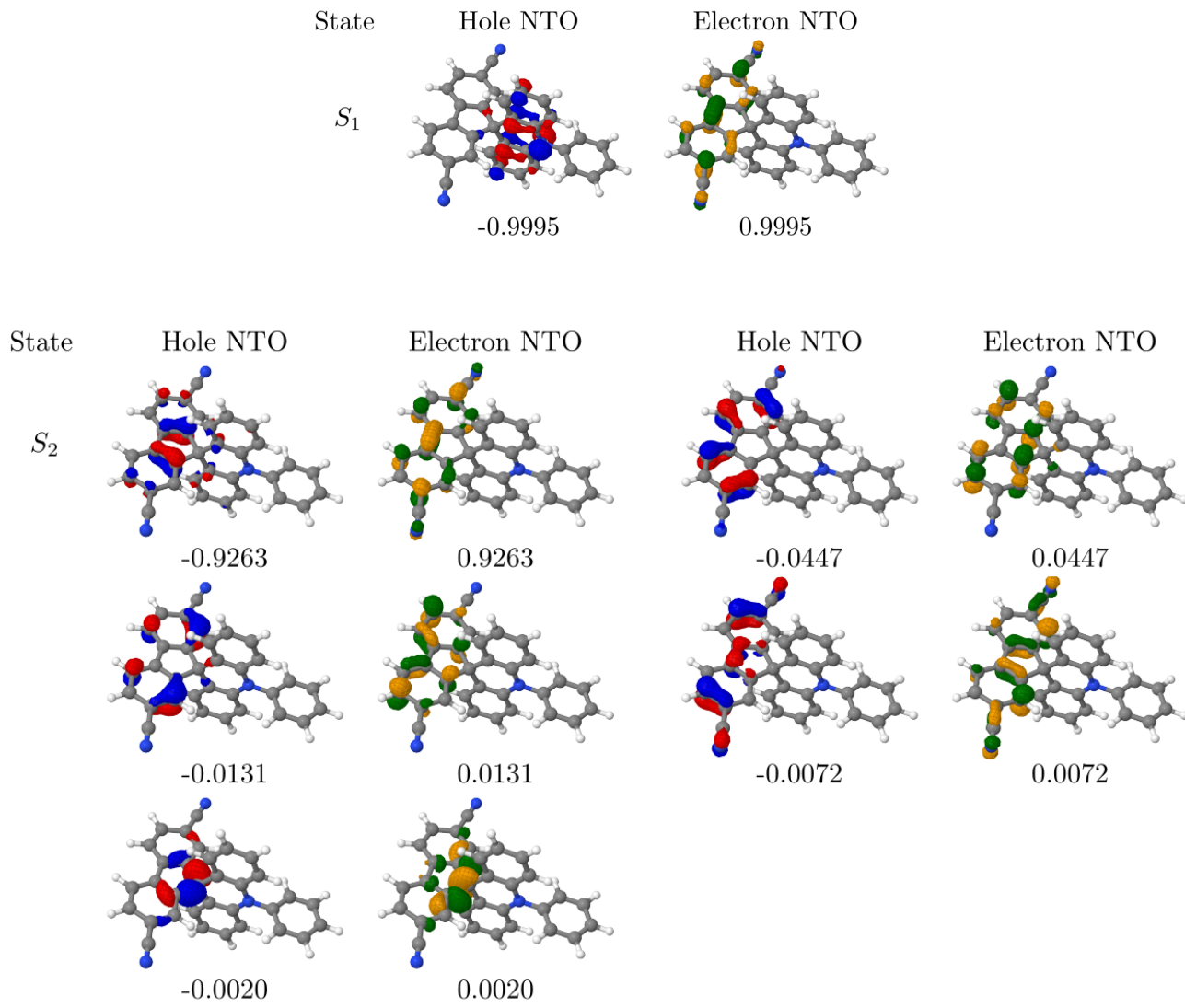


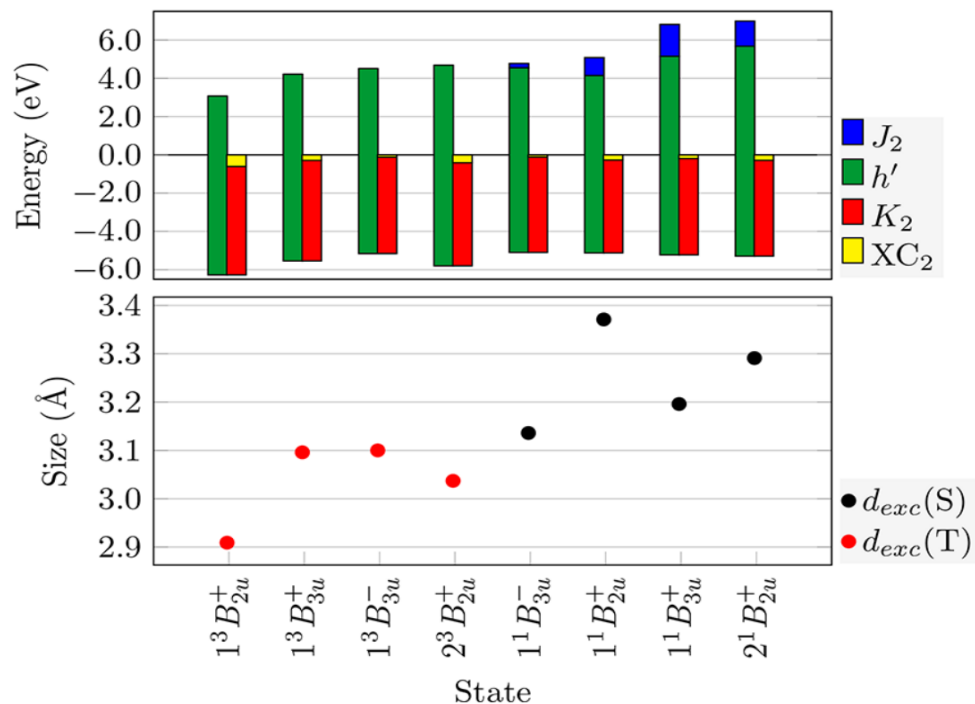
## Why the lowest excited state is not always the $H \rightarrow L$ Transition?

- The Coulomb attraction ( $K_2$ ) stabilizes locally excited states, especially triplets, while the exchange repulsion ( $J_2$ ) destabilizes locally excited singlet states.
- In **ACRFLCN**, the lowest triplet state,  $T_1$ , originates from the  $H-1 \rightarrow L$  transition because the stronger  $K_2$  stabilizes this configuration below the  $H \rightarrow L$  transition.
- For **naphthalene**, the lowest  $S_1$  arises from a linear combination of  $H-1 \rightarrow L$  and  $H \rightarrow L+1$  transitions because  $J_2$  destabilize the  $H \rightarrow L$  state and shifts it to higher energy.









### Excited energy components

$$T_1(1^3B_{2u}^+) = 3.07 \text{ eV} = 9.35 + (-5.67) + (-0.61)$$

$$T_2(1^3B_{3u}^+) = 4.21 \text{ eV} = 9.76 + (-5.26) + (-0.29)$$

$$T_3(1^3B_{3u}^-) = 4.51 \text{ eV} = 9.67 + (-5.05) + (-0.11)$$

$$T_4(2^3B_{2u}^+) = 4.68 \text{ eV} = 10.48 + (-5.39) + (-0.41)$$

$$S_1(1^1B_{3u}^-) = 4.78 \text{ eV} = 9.65 + 0.23 + (-4.99) + (-0.11)$$

$$S_2(1^1B_{2u}^+) = 5.09 \text{ eV} = 9.27 + 0.94 + (-4.85) + (-0.27)$$

$$S_3(1^1B_{3u}^+) = 6.82 \text{ eV} = 10.37 + 1.67 + (-5.03) + (-0.19)$$

$$S_4(2^1B_{2u}^+) = 6.99 \text{ eV} = 10.97 + 1.31 + (-5.01) + (-0.28)$$

### Main Configuration

H → L

|H-1 → L| + |H → L+1|

|H → L+1| - |H-1 → L|

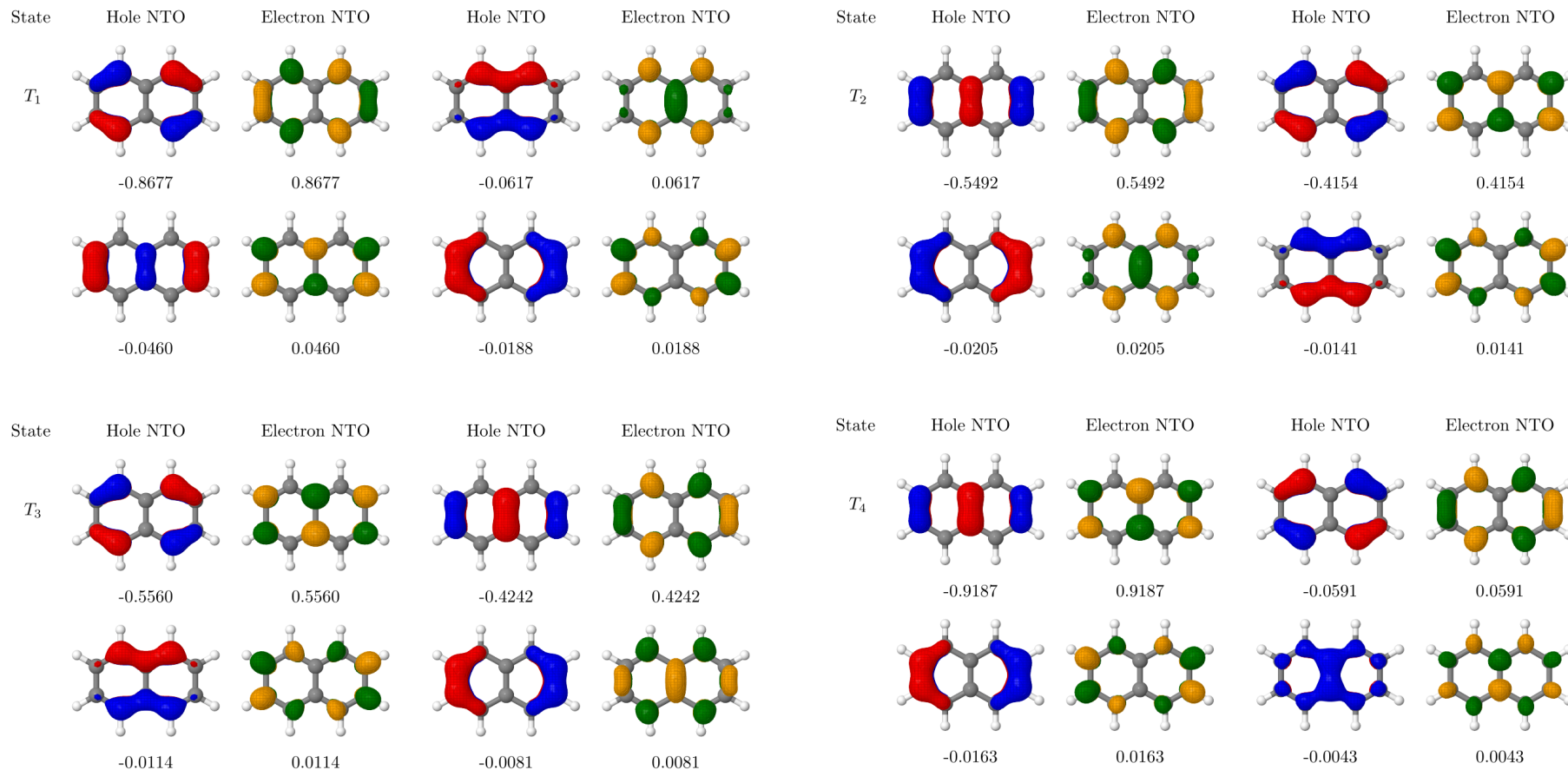
H-1 → L+1

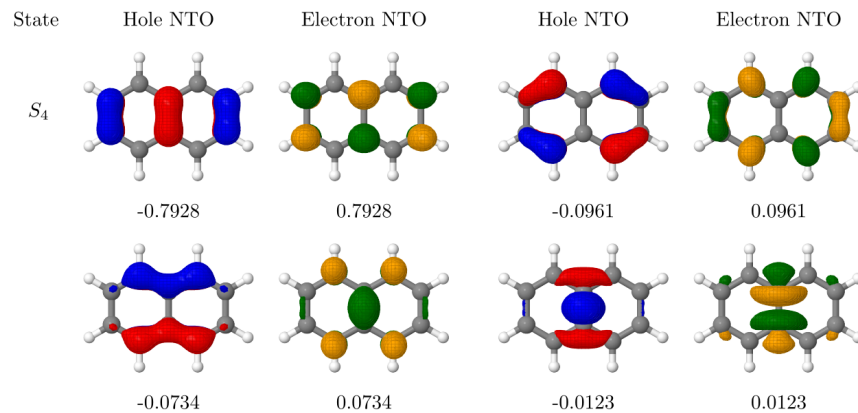
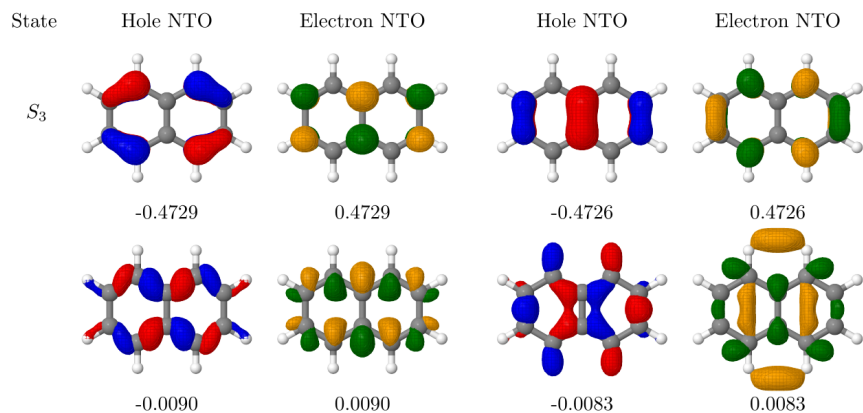
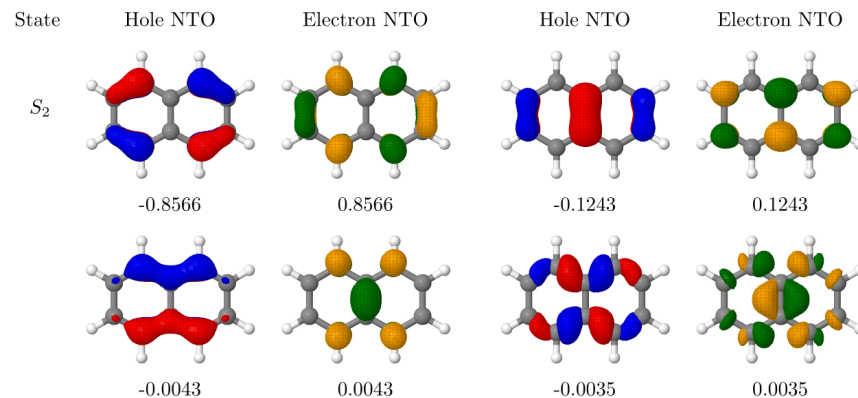
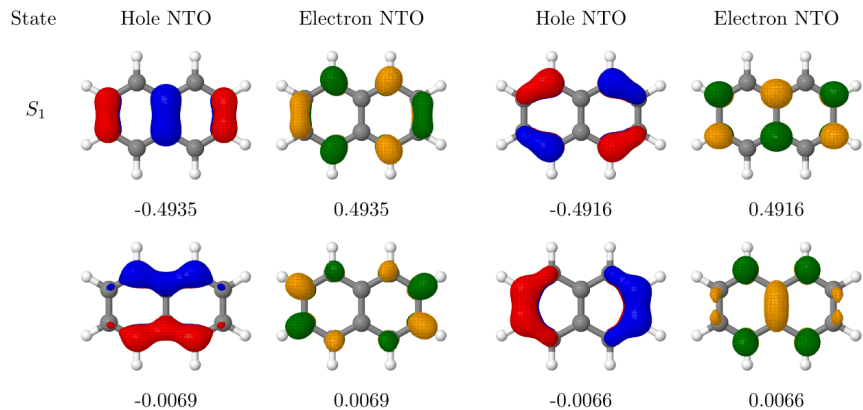
|H → L+1| - |H-1 → L|

H → L

|H-1 → L| + |H → L+1|

H-1 → L+1





## Electrostatic potential (ESP)

$$\begin{aligned}\iint \frac{\rho_A(r_1)\rho_B(r_2)}{r_{12}} dr_1 dr_2 &= \int \rho_A(r_1) \left( \int \frac{\rho_B(r_2)}{r_{12}} dr_2 \right) dr_1 \\ &= \int \rho_A(r_1) \hat{V} \rho_B(r_1) dr_1 = \langle \rho_A | \hat{V} | \rho_B \rangle.\end{aligned}$$

## Natural transition orbital (NTO)

The 1-electron transition density matrix (1TDM) is defined as

$$\gamma_{0I}(r_h, r_e) = \sum_{pq} \gamma_{pq}^{0I} \phi_p(r_h) \phi_q(r_e)$$

Contributed MO

where,

$$\gamma_{pq}^{0I} = \langle \Psi_0 | \hat{a}_p^\dagger \hat{a}_q | \Psi_I \rangle$$

Creation and annihilation operator of MO

Applying the singular value decomposition,

$$\gamma_{0I}(r_h, r_e) = \sum_t \sqrt{\lambda_t} \psi_t^h(r_h) \psi_t^e(r_e)$$

Amplitude

NTOs representing the hole and electron